

ponents. Their amplitude is small and the period is larger than the entry length. These disturbances are not found in the absence of the Hall and ion slip currents. This increases the entry length. This conclusion agrees with the results of Saric and Touryan.⁶

Numerically, the entry length has been calculated on the basis of 5% deviations in the fully developed centerline velocity. The entry lengths and K_{fd} , the fully developed values of the correction term K due to entrance effects in the pressure distribution are shown in Table 1.

The entry length for the case of no Hall and ion slip currents with Hartmann number $H_a = 10$ was found to be $0.0143 R_e h$ by Chen and Chen.⁴ With the Hall parameter $\beta_e = 2.0$ it is $0.0207 R_e h$ and with $\beta_e = 2.0$ and ion slip parameter $\beta_i = 0.5$, it is found to be $0.0216 R_e h$. As Hartmann number increases the entrance length decreases sharply.

The pressure distribution in the entrance region is shown in Fig. 1 for different values of H_a , β_e , and β_i . It shows small curvature in the inlet region. For the downstream region it becomes constant. The solutions presented herein may be useful for the investigation of the stability of a developing MHD flow with Hall and ion slip currents.

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Correlation for Gasdynamic Laser Gain

James I. McManus* and John D. Anderson Jr.†
University of Maryland, College Park, Md.

Introduction

GASDYNAMIC lasers (GDL's) have spearheaded a breakthrough in high energy laser technology. These devices are essentially supersonic wind tunnels which create a lasing medium by rapid expansion of a vibrationally excited

molecular gas. Under suitable conditions, laser power can be extracted from the supersonic stream. The physical processes, practical importance, experimental data, and theoretical analyses associated with GDL's are described in a recent book.¹

Theoretical prediction of gasdynamic laser performance generally requires a sophisticated nonequilibrium nozzle flow computer program, such as Ref. 2. Moreover, GDL performance is a function of numerous variables, such as nozzle shape and size, reservoir gas temperature and pressure, mixture ratio, etc. For these reasons, an analysis of gasdynamic laser characteristics is usually restricted to those individuals and organizations which have considerable computational capability. Therefore, there is a need for engineering correlations which allow quick and easy hand calculations of GDL performance without a gross compromise in accuracy. The purpose of the present Note is to provide such an engineering correlation. In particular, a formula is provided for the calculation of peak small-signal gain for $\text{CO}_2\text{-N}_2\text{-H}_2\text{O}$ gasdynamic lasers. Such a mixture is common in practical GDL field and laboratory devices.

The Correlation

Small-signal gain, G_o , is a measure of the amplifying property of the laser medium such that $dI_o = G_o I_o dx$, where I_o is the laser beam intensity, and dI_o is the increase in intensity over a distance dx . Essentially, G_o is a negative absorption coefficient. In gasdynamic lasers, G_o varies with distance in the flow direction, first increasing, reaching a peak, then decreasing downstream.¹ In the present Note, the peak value of G_o is correlated as the following functional variation,

$$G_o = f(P_o, T_o, A_e/A^*, h^*, X_{N_2}, X_{\text{CO}_2}, X_{\text{H}_2\text{O}}) \quad (1)$$

where P_o and T_o are the reservoir gas pressure and temperature, respectively, A_e/A^* is the nozzle exit-to-throat area ratio, h^* is the throat height, and X_i is the mole fraction of species i in the mixture.

Using the computer program described in Ref. 2, with updated kinetic rates as given in Ref. 3, a large number of parametric variations were obtained. The parameters were those of Eq. (1), with the exception that P_o and h^* were grouped as the product $P_o h^*$, based on binary scaling, and $X_{\text{CO}_2} + X_{\text{N}_2} + X_{\text{H}_2\text{O}} = 1$. This large bulk of numerical data was then systematically correlated, using a combination of chi-square and Gaussian functions, aided by least-squares fits. The details are given in Ref. 4. The resulting engineering formula, although rather long, can easily be evaluated on a slide rule or pocket calculator. The formula is

$$G_o = K_o - G[(1-Q1) + (1-Q2) + (1-Q3)] - Q4 \quad (2)$$

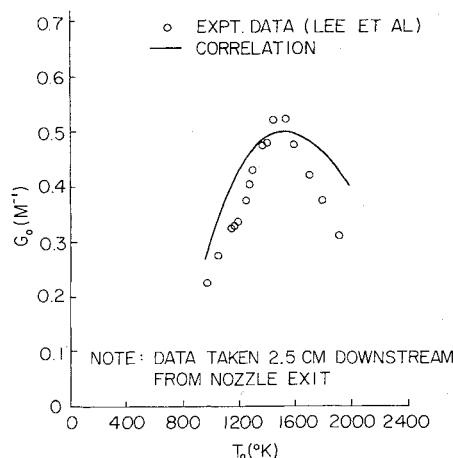


Fig. 1 Comparison of Eq. (2) with data of Ref. 5. $h^* = 0.1$ cm, $A_e/A^* = 15$ ($M = 4$), $P_o = 9$ atm, $\% \text{H}_2\text{O} = 2.4$, $\% \text{CO}_2 = 6.3$.

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*Graduate Student, Department of Aerospace Engineering. Member AIAA.

†Professor and Chairman, Department of Aerospace Engineering. Associate Fellow AIAA.

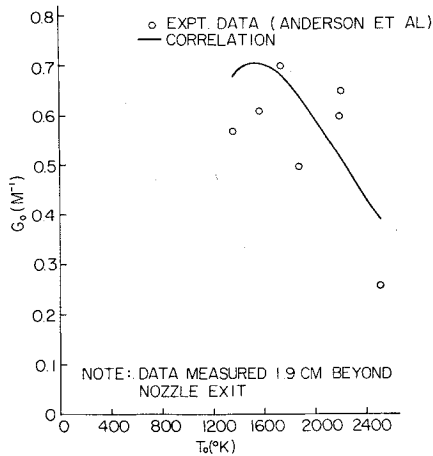


Fig. 2 Comparison of Eq. (2) with data of Ref. 6. $h^* = 0.1$ cm, $A_e/A^* = 20$, $P_o = 10.23$ atm, $\%H_2O = 1.2$, $\%CO_2 = 10.3$.

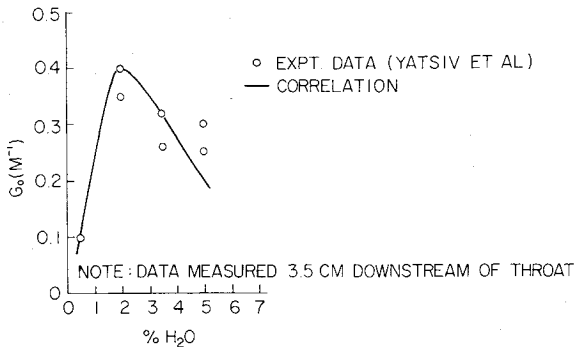


Fig. 3 Comparison of Eq. (2) with data of Ref. 7. $h^* = 0.13$ cm, $A_e/A^* = 15$ (wedge nozzle with 15° half-angle), $P_o = 7.5$ atm, $T_o = 1650$ K, $\%CO_2 = 20$.

where

$$K_O = 1.484 e^{-(A_e/A^* - 53.037)^2 / 1532.3}$$

$$G = 0.875 K_O$$

$$Q1 = 0.025 T_X^{2.673} e^{-0.2475 T_X}$$

$$Q2 = 2.75 \sin(C_X) e^{-C_X / 1.171}$$

$$Q3 = 0.8 (H_2O^1 + 1.0)^{0.65} e^{-0.245 H_2O^1}$$

$$Q4 = 0.1655 \exp \left[\frac{-(A_e/A^* - 39.15)^2}{176.72} \right] (P_o h^*)^{3/2}$$

$$T_X = [T_o / (1.0 - .105 AF) - 470.] / 100.0$$

$$C_X = (\%CO_2 \pi) / (40. + 11. AF)$$

$$H_2O^1 = (\%H_2O - CA) (\%CO_2 / 24 - CA / 3.3) (1.05 - 0.35 AF) + 1.65$$

$$CA = 2.25 - 2.0 e^{-\%CO_2 / 10.0}$$

$$AF = (A_e/A^* - 19.75) / 30.25$$

$$T_o = \text{stagnation temperature (K)}$$

$$\%CO_2 = \text{mole fraction of } CO_2$$

$$\%H_2O = \text{mole fraction of } H_2O$$

$$P_o = \text{stagnation pressure (atm)}$$

$$h^* = \text{throat height (cm)}$$

$$A_e/A^* = \text{nozzle exit area to throat area ratio}$$

Equation (2) is valid only for minimum-length nozzles, $A_e/A^* < 53$, $0.5 < \%H_2O < 30$, $\%CO_2 < 30$ and $P_o h^* < 3.0$. However, these ranges include almost all cases of practical interest.

Results

As shown in Ref. 4, Eq. (2) correlates the numerical data of Refs. 2 and 3 within 1 or 2% for most of the runs made.

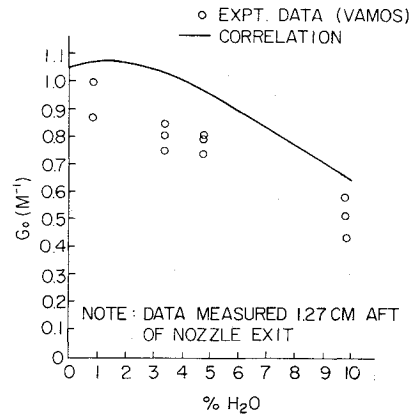


Fig. 4 Comparison of Eq. (2) with data of Ref. 8. $h^* = 0.0356$ cm, $A_e/A^* = 50$, $P_o = 37.5$ atm, $T_o = 1800$ K, $\%CO_2 = 7$.

However, the real test of the formula is its application to experimental conditions. Figures 1-4 show such a comparison, where Eq. (2) is matched with the experimental data of Lee et al.,⁵ Anderson et al.,⁶ Yatsiv et al.,⁷ and Vamos,⁸ respectively. The data are chosen because they encompass a wide range of area ratio, temperature, and mixture ratio. In all cases, the comparison is reasonable. However, emphasis is made that the correlation is for peak G_o , whereas the experimental data are obtained at fixed locations in the laser cavity which do not necessarily correspond to the location of peak gain. Also, Eq. (2), being based on the numerical data of Refs. 2 and 3, is therefore subject to the same shortcomings of the theoretical analysis and kinetic rates. Nevertheless, the correlation appears to be reasonably accurate. This is particularly remarkable considering the complicated dependency of G_o on the parameters given in Eq. (1).

Conclusions

A closed-form correlation is given for the peak small-signal gain in CO_2 - N_2 - H_2O gasdynamic lasers. This is the first step towards quick and easy "back-of-the-envelope" estimates of GDL performance. The correlation agrees reasonably well with experimental data. In view of the fact that the correlations encompasses a large number of GDL parameters over wide ranges, such agreement is rather remarkable.

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